

VARIATIONAL NEURAL NETWORK ANSATZ FOR STEADY-STATES IN OPEN QUANTUM SYSTEMS

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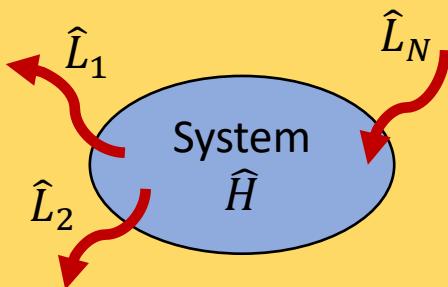
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ARXIV:1902.10104 [quant-ph]

OPEN QUANTUM SYSTEMS

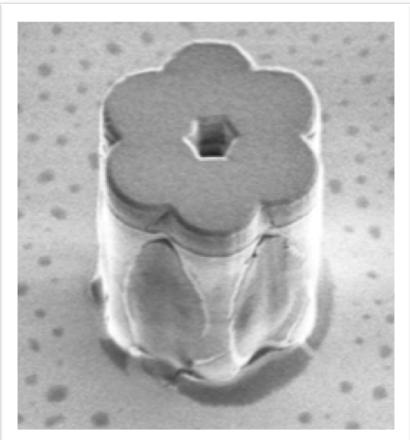
Describe the interaction of a quantum system H with its environment

Reservoir



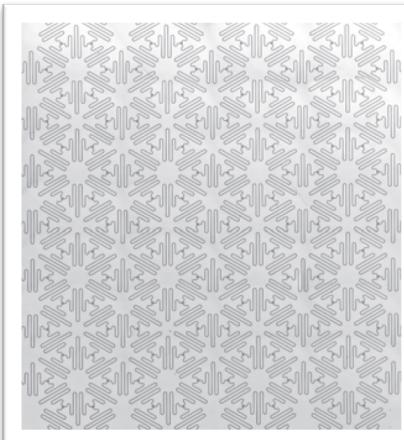
- Out of Equilibrium phenomena
- Reservoir Engineering
- *Dissipative* phase transitions

Semiconductor Micropillars



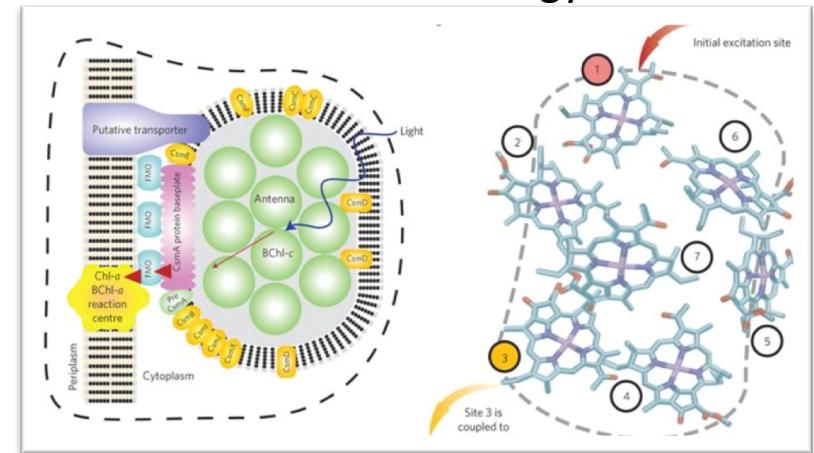
N. Carlon Zambon et Al.
arXiv:1812.06163

Superconducting circuits



A. A. Houck et Al.
Nat Phys 8, 292 (2012)

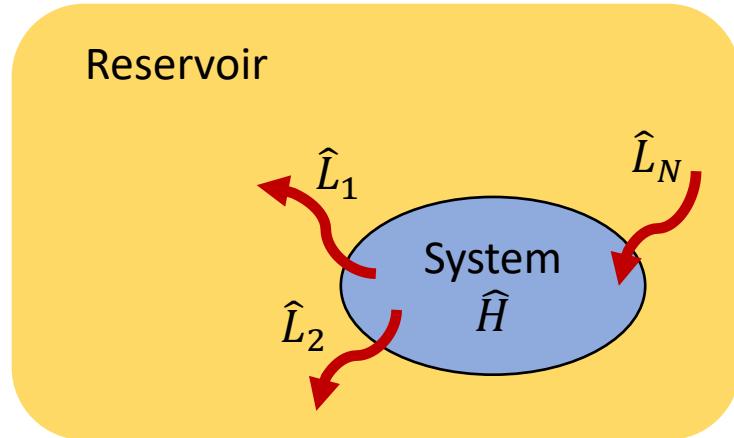
Quantum Biology



N. Lambert et Al.
Nat Phys 9, 10 (2013)

SCHROEDINGER'S EQUATION → LINDBLAD MASTER EQ.

Describe the interaction of a quantum system H with its environment



THE STATE

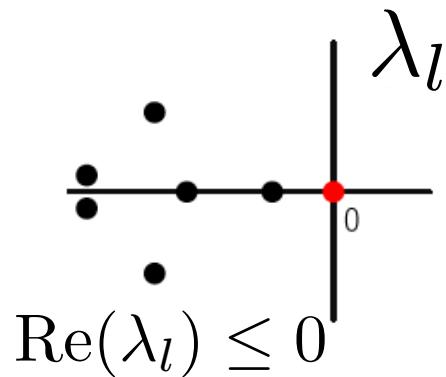
$$|\psi\rangle \in \mathcal{H}_{\text{sys}} \otimes \mathcal{H}_{\text{E}} \quad \Rightarrow \quad \hat{\rho} = \text{Tr}_{\text{E}}[|\psi\rangle \langle \psi|]$$

THE MASTER EQUATION

$$\frac{\partial \hat{\rho}}{\partial t} = -i[\hat{H}, \hat{\rho}] + \sum_{j=1}^{N_{\text{channels}}} \left(\hat{L}_j \hat{\rho} \hat{L}_j^\dagger - \frac{1}{2} \{ \hat{L}_j^\dagger \hat{L}_j, \hat{\rho} \} \right)$$

COHERENT EVOLUTION

INCOHERENT EVOLUTION



REWRITTEN AS A LINEAR SUPER-OPERATORIAL EQUATION

$$\partial_t \hat{\rho} = \mathcal{L} \hat{\rho}$$

WE ARE USUALLY INTERESTED IN THE STEADY-STATE

$$\hat{\rho}_{ss} = \lim_{t \rightarrow \infty} e^{\mathcal{L}t} \hat{\rho}$$

$$\begin{aligned} \mathcal{L} \hat{\rho}_{ss} &= 0 \\ \partial_t \hat{\rho}_{ss} &= 0 \end{aligned}$$

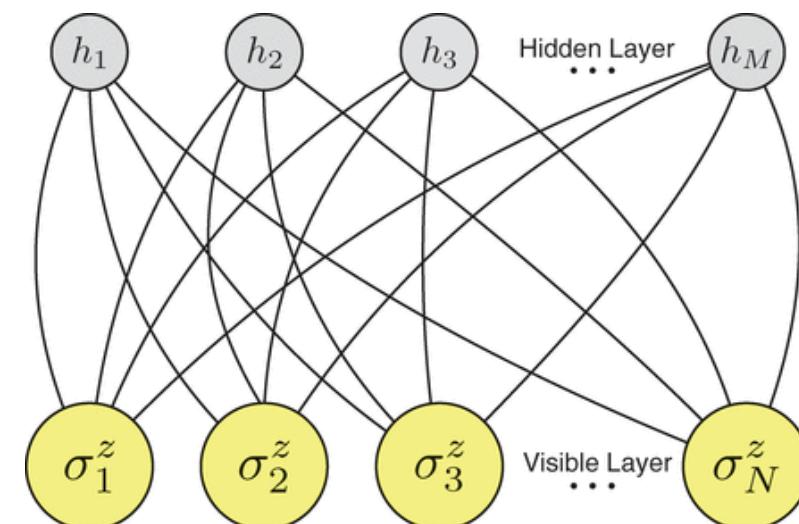
THE DIMENSIONAL CURSE

FOR N LATTICES OF N SITES, THE HILBERT SPACE HAS AN EXPONENTIALLY BIG DIMENSION $\dim \mathcal{H} \approx l^N$

NUMERICAL TECHNIQUES TRY TO REDUCE THIS SIZE WITH PHYSICAL INSIGHT (EG: MPS/MPOs, CLUSTER MEAN FIELD..)

NEURAL NETWORKS ARE LOW-DIMENSIONAL APPROXIMATIONS OF HIGH-DIMENSIONAL FUNCTIONS...

$$|\psi\rangle = \begin{pmatrix} \psi(\uparrow\uparrow \dots \uparrow) \\ \psi(\downarrow\uparrow \dots \uparrow) \\ \psi(\uparrow\downarrow \dots \uparrow) \\ \psi(\downarrow\downarrow \dots \uparrow) \\ \vdots \\ \psi(\downarrow\downarrow \dots \downarrow) \end{pmatrix}$$



THE DIMENSIONAL CURSE

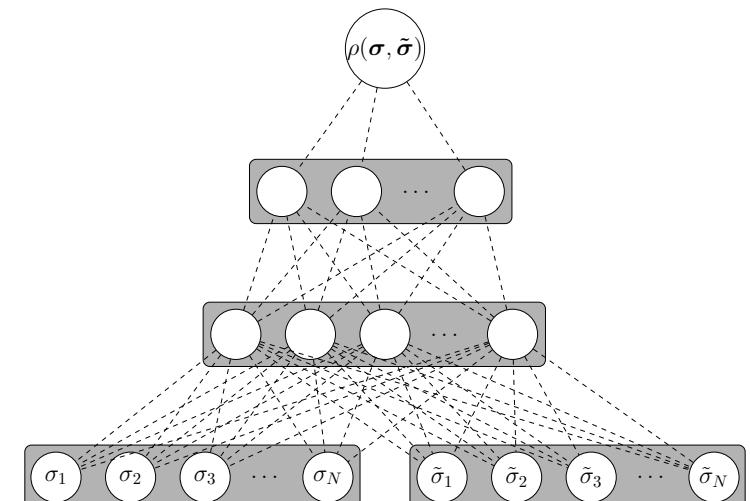
FOR N SPINS, THE HILBERT SPACE HAS AN EXPONENTIALLY BIG DIMENSION $\dim \hat{\rho} = \dim \mathcal{H} \otimes \mathcal{H} = 2^{2N}$

$$\hat{\rho} = \sum_{\substack{\sigma_1, \dots, \sigma_N, \\ \tilde{\sigma}_1, \dots, \tilde{\sigma}_N}} \underbrace{\rho(\sigma_1, \dots, \sigma_N, \tilde{\sigma}_1, \dots, \tilde{\sigma}_N)}_{\text{HIGH-DIMENSIONAL FUNCTION}} |\sigma_1, \dots, \sigma_N\rangle \langle \tilde{\sigma}_1, \dots, \tilde{\sigma}_N|$$

HIGH-DIMENSIONAL FUNCTION

4 ARTICLES CAME OUT IN A WEEK WITH SIMILAR PROPOSALS

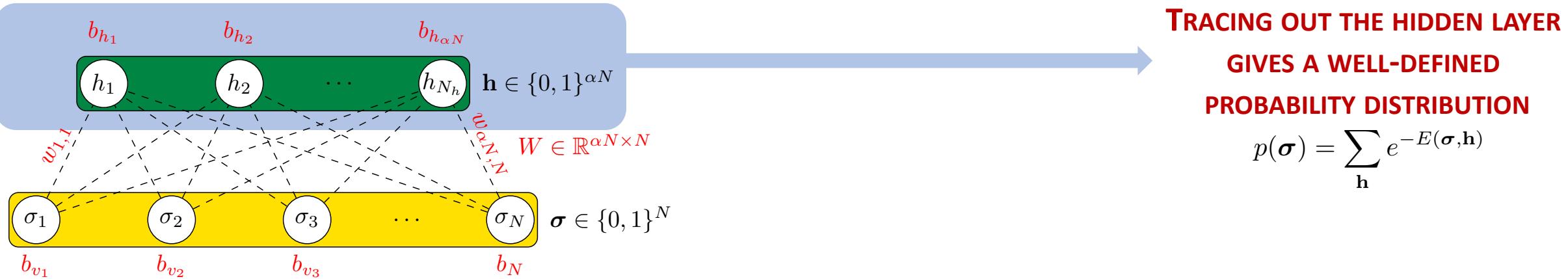
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|---|------------------|
| F.V, A. BIELLA, N. REGNAULT AND C.CIUTI | ARXIV:1902.10104 |
| M.J. HARTMANN AND G.CARLEO | ARXIV:1902.05131 |
| A. NAGY AND V. SAVONA | ARXIV:1902.09483 |
| N. YOSHIOKA AND R. HAMAZAKI | ARXIV:1902.07006 |



BUT $\hat{\rho}$ IS POSITIVE-SEMIDEFINITE AND HERMITIAN. WE WANT TO ENFORCE THOSE PROPERTIES.

NEURAL DENSITY MATRIX

Drawing inspiration from RBMs (Restricted Boltzmann Machines)...

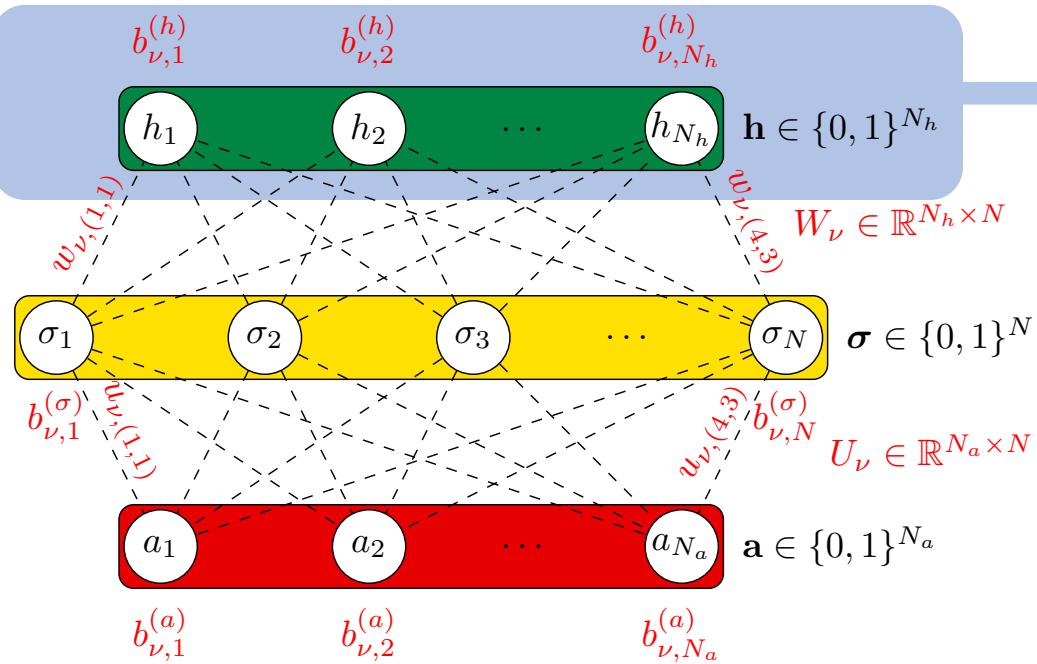


TRACING OUT THE HIDDEN LAYER
GIVES A WELL-DEFINED
PROBABILITY DISTRIBUTION

$$p(\boldsymbol{\sigma}) = \sum_{\mathbf{h}} e^{-E(\boldsymbol{\sigma}, \mathbf{h})}$$

NEURAL DENSITY MATRIX

Drawing inspiration from RBMs (Restricted Boltzmann Machines)...



**TRACING OUT THE HIDDEN LAYER
GIVES A WELL-DEFINED
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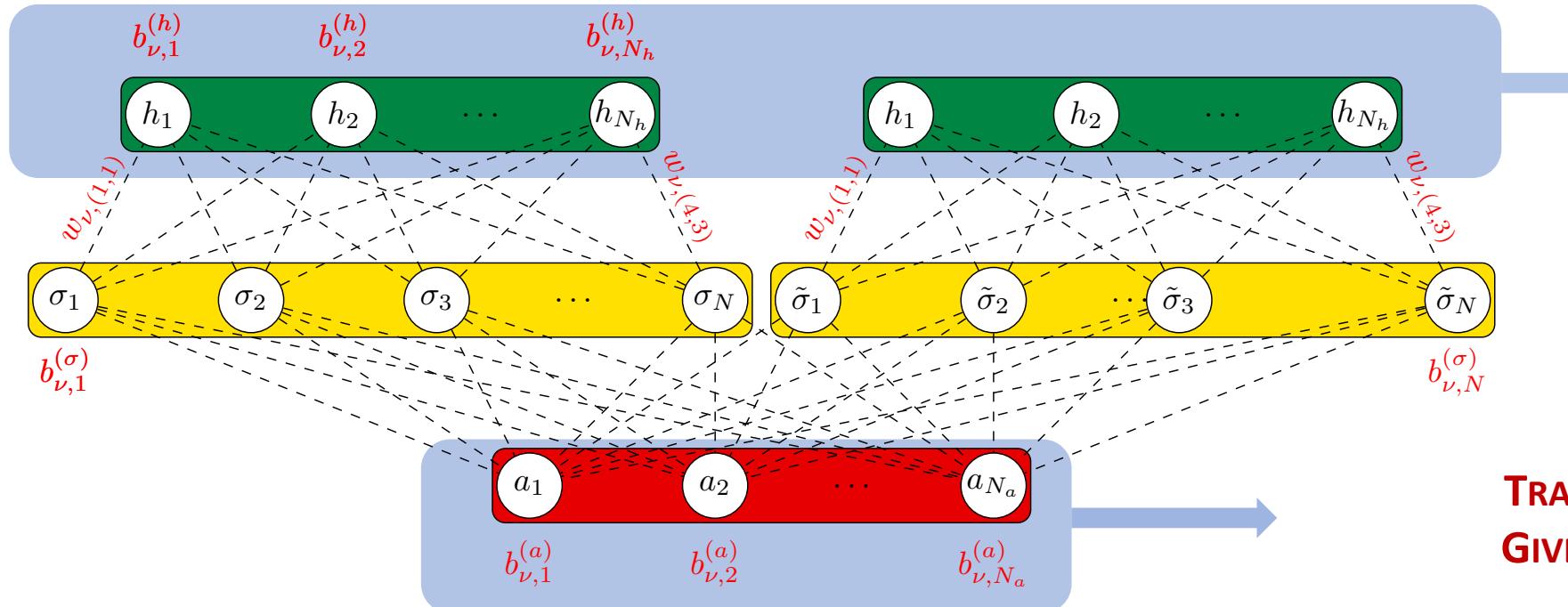
$$p(\boldsymbol{\sigma}) = \sum_{\mathbf{h}} e^{-E(\boldsymbol{\sigma}, \mathbf{h})}$$

DOUBLE THE «HIDDEN» VARIABLES

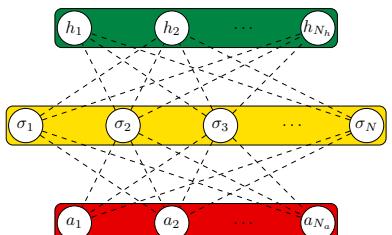
$$p(\boldsymbol{\sigma}, \mathbf{a}) = \sum_{\mathbf{h}} e^{-E(\boldsymbol{\sigma}, \mathbf{a}, \mathbf{h})}$$

NEURAL DENSITY MATRIX

Drawing inspiration from RBMs (Restricted Boltzmann Machines)...



Formally equivalent to defining an RBM in the bigger space, summing over the hidden space and then



$$\psi(\vec{\sigma}, \vec{a}) \in \mathcal{H}_s \otimes \mathcal{H}_a \quad \Rightarrow \quad \rho(\vec{\sigma}, \vec{\sigma}') = \sum_{a \in \mathcal{H}_a} [\psi(\vec{\sigma}, \vec{a}) \psi^*(\vec{\sigma}', \vec{a})]$$

Can be generalized to FFNN!

OUT OF EQUILIBRIUM: FINDING THE STEADY STATE

We have a variational ansatz (where v are the variational parameters)

$$\rho_v(\sigma, \sigma')$$

I want to find the steady state

$$\frac{d\rho_{ss}}{dt} = \mathcal{L}\rho_{ss} = 0$$

We need a variational principle: I define the 'Cost Function'

$$\mathcal{C}(v) = \frac{\|d\hat{\rho}_v/dt\|_2^2}{\|\hat{\rho}_v\|_2^2} = \frac{\text{Tr}[\hat{\rho}_v^\dagger \mathcal{L}^\dagger \mathcal{L} \hat{\rho}_v]}{\text{Tr}[\hat{\rho}_v^\dagger \hat{\rho}_v]}$$

Which has the nice properties:

$$\mathcal{C}(v_{ss}) = 0 \iff \hat{\rho}_{v_{ss}} = \hat{\rho}_{ss}$$

$$\mathcal{C}(v) \geq 0$$

DIFFERENT COST FUNCTIONS ARE POSSIBLE!
SEE POSTER OF A. NAGY FOR A DIFFERENT ONE



OUT OF EQUILIBRIUM: FINDING THE STEADY STATE III

We recasted the problem of finding the steady state ρ_{ss} to a minimization problem for $\mathcal{C}(\mathbf{v})$

So we could perform a gradient-descent like minimization: $\mathbf{v} \rightarrow \mathbf{v} - \nabla_{\mathbf{v}} \mathcal{C}(\mathbf{v})$

But there's a problem! Local Minimas! → **NATURAL GRADIENT DESCENT**

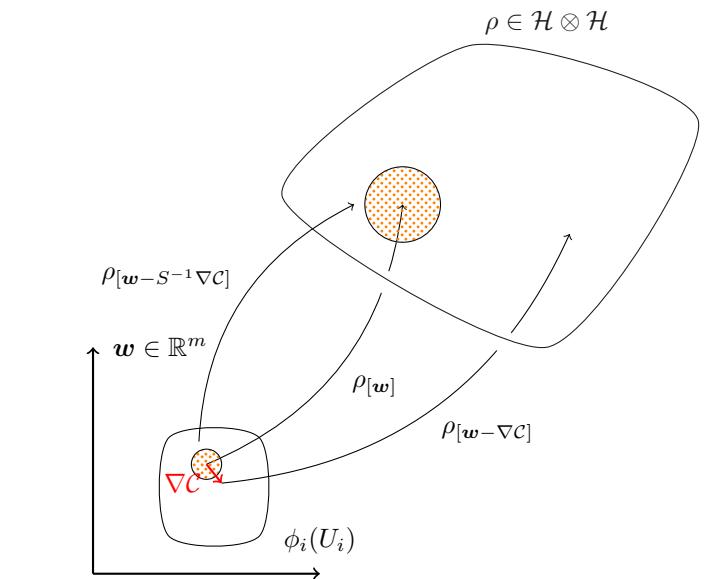
But there's a problem! Computing the Cost Fun → **SAMPLING**

$$\mathcal{C}(\mathbf{v}) = \frac{\text{Tr}[\hat{\rho}_{\mathbf{v}}^\dagger \mathcal{L}^\dagger \mathcal{L} \hat{\rho}_{\mathbf{v}}]}{\text{Tr}[\hat{\rho}_{\mathbf{v}}^\dagger \hat{\rho}_{\mathbf{v}}]} = \sum_{\sigma, \tilde{\sigma}} p_{\mathbf{v}}(\sigma, \tilde{\sigma}) \mathcal{C}^{\text{Loc}}(\mathbf{v}, \sigma, \tilde{\sigma}),$$

$$p_{\mathbf{v}}(\sigma, \tilde{\sigma}) = |\rho_{\mathbf{v}}(\sigma, \tilde{\sigma})|^2 / \sum_{\sigma, \tilde{\sigma}} \rho_{\mathbf{v}}(\sigma, \tilde{\sigma})$$

Exponentially
Many elements

This is a probability distribution → I can sample it! (and the gradient)



$$\mathcal{C}^{\text{Loc}}(\mathbf{v}, \sigma, \tilde{\sigma}) = \sum_{\sigma', \tilde{\sigma}'} \frac{\rho_{\mathbf{v}}(\sigma', \tilde{\sigma}')}{\rho_{\mathbf{v}}(\sigma, \tilde{\sigma})} (\mathcal{L}^\dagger \mathcal{L})(\sigma, \tilde{\sigma}; \sigma', \tilde{\sigma}').$$

RESULTS FOR DRIVEN-DISSIPATIVE QUANTUM ISING MODEL

The hamiltonian is

$$H = \frac{V}{4} \sum_{\langle j, l \rangle} \hat{\sigma}_j^z \hat{\sigma}_l^z + \frac{g}{2} \sum_j \hat{\sigma}_j^x$$

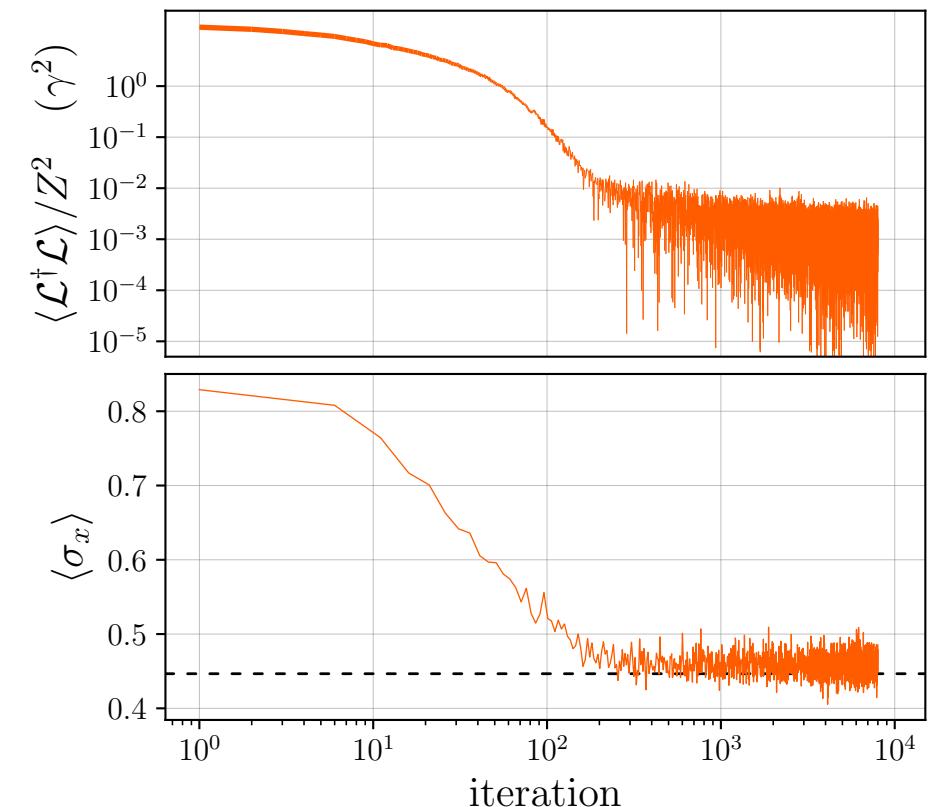
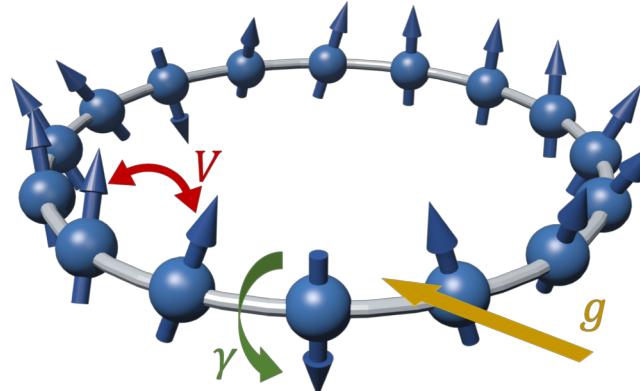
With de-excitation jump operators $L_j = \sigma_j^-$

We study the magnetizations $m^{(\alpha)} = \frac{1}{N} \sum_i \sigma_i^{(\alpha)}$

For a chain of N=16 sites (PBC)

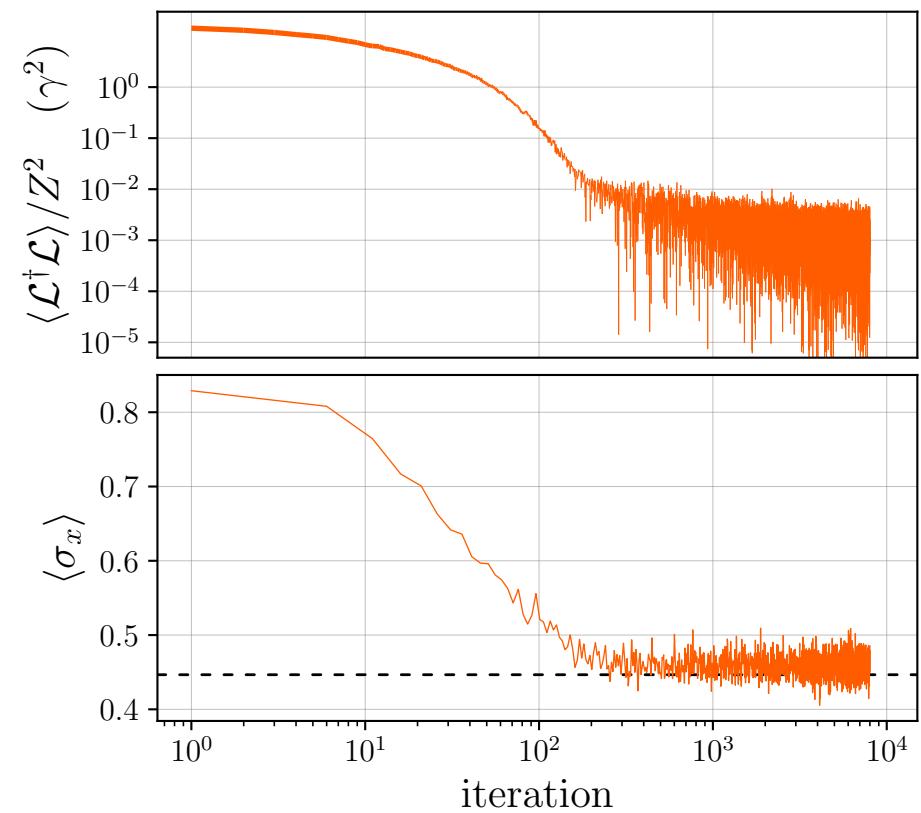
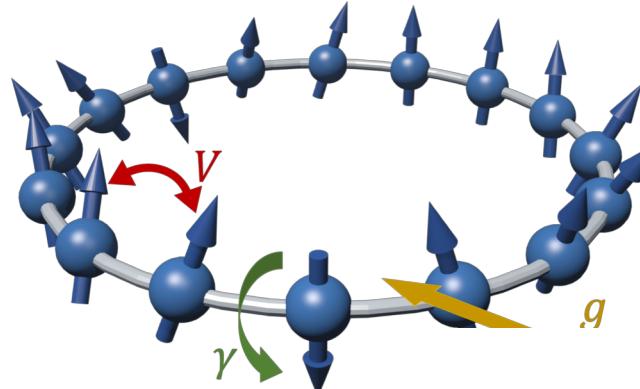
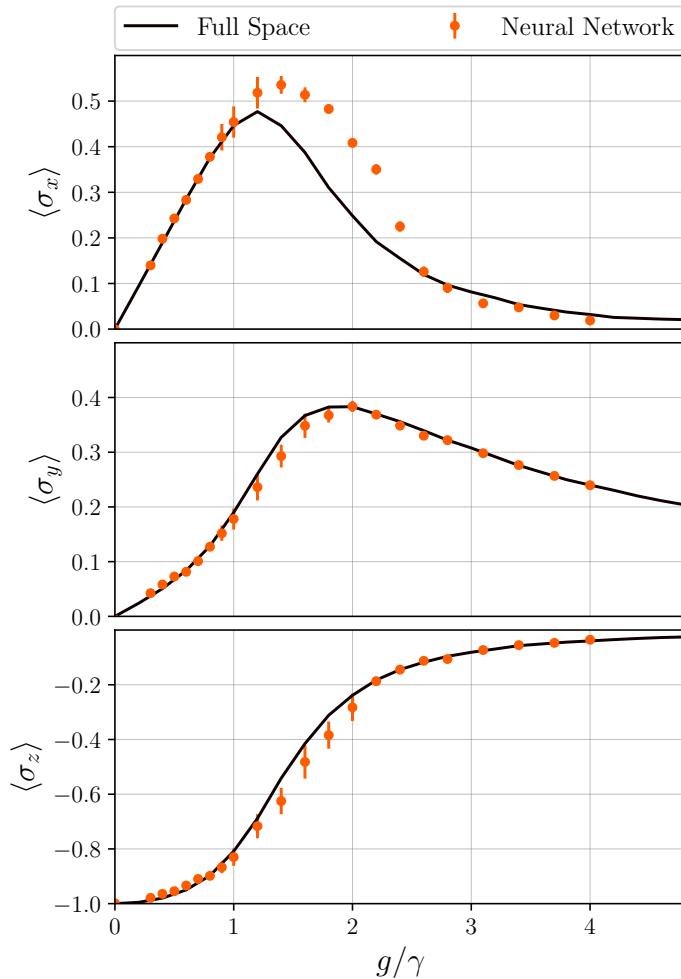
Observables are sampled like:

$$\langle \hat{m}^{(\alpha)} \rangle = \frac{\text{Tr}[\hat{\rho} \hat{m}^{(\alpha)}]}{\text{Tr}[\hat{\rho}]} = \sum_{\sigma} p_{\sigma}^{\text{obs}}(\sigma) \sum_{\tilde{\sigma}} \frac{\rho_{\sigma}(\sigma, \tilde{\sigma}) m^{(\alpha)}(\tilde{\sigma}, \sigma)}{\rho_{\sigma}(\sigma, \sigma)},$$



RESULTS FOR DRIVEN-DISSIPATIVE QUANTUM ISING MODEL II

A full scan in the transverse field gives:



CONCLUSIONS

- Density Matrices can be approximated with Neural Networks
- Variational Monte Carlo can be remapped to machine learning procedures
- We can solve the exponential growth problem

Perspectives:

- Try different topologies, layers to enforce simmetries...
- Test different cost functions
- [Also investigate Neural Network encodings for Digital Quantum Algorithm applications]

AKNOWLEDGEMENTS

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«Variational neural Network ansatz for steady states in Open Quantum Systems»

ArXiv:1902.10104



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QUESTIONS?

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CONTACT ME

QUESTIONS?

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CODE?

Will soon be available on GitHub. If you want to have a look now drop me an email.

COLLABORATE?

I am finalizing a framework written in Julia to perform MonteCarlo simulations using Neural Network ansatzes. If you would like to collaborate and/or help, I would love an hand.